Zeitschrift: Helvetica Physica Acta

Band: 46 (1973)

Heft: 5

Artikel: Total spontaneous fission half-lives, kinetic energy and mass yield

spectra of 250Cm, 254Cf and 258Fm

Autor: Hooshyar, M.A. / Bary Malik, F.

DOI: https://doi.org/10.5169/seals-114506

Nutzungsbedingungen

Die ETH-Bibliothek ist die Anbieterin der digitalisierten Zeitschriften. Sie besitzt keine Urheberrechte an den Zeitschriften und ist nicht verantwortlich für deren Inhalte. Die Rechte liegen in der Regel bei den Herausgebern beziehungsweise den externen Rechteinhabern. Siehe Rechtliche Hinweise.

Conditions d'utilisation

L'ETH Library est le fournisseur des revues numérisées. Elle ne détient aucun droit d'auteur sur les revues et n'est pas responsable de leur contenu. En règle générale, les droits sont détenus par les éditeurs ou les détenteurs de droits externes. Voir Informations légales.

Terms of use

The ETH Library is the provider of the digitised journals. It does not own any copyrights to the journals and is not responsible for their content. The rights usually lie with the publishers or the external rights holders. See Legal notice.

Download PDF: 02.04.2025

ETH-Bibliothek Zürich, E-Periodica, https://www.e-periodica.ch

Total Spontaneous Fission Half-Lives, Kinetic Energy and Mass Yield Spectra of ²⁵⁰Cm, ²⁵⁴Cf and ²⁵⁸Fm

by M. A. Hooshyar

Physics Department, Arya-Meher University of Technology, Tehren, Iran

and F. Bary Malik¹)

Institut de Physique, Université de Neuchâtel, Switzerland

(3. VII. 73)

Abstract. i) The absolute spontaneous fission half-lives of ²⁵⁰Cm, ²⁵⁴Cf and ²⁵⁸Fm, ii) the preference of ²⁵⁸Fm to decay in symmetric modes, and iii) trends of half-lives of even Curium and Californium isotopes are accounted for within the context of the recent theory.

Among isotopes of Curium, Californium and Fermium, 250 Cm, 254 Cf and 258 Fm have the shortest half-lives against spontaneous fission. It has not yet been possible, theoretically, to account for this fact as well as their absolute total spontaneous fission half-lives. In addition, recent measurements indicate that dominant decay modes of 258 Fm are near the symmetric mass splitting [1] and 258 Fm has an extremely short spontaneous fission half-life [2] of $(3.8 \pm 0.6)~10^{-4}$ sec. Thus 258 Fm is one of the first detected isotopes decaying mostly to symmetric decay modes. Since this is also the heaviest isotope whose spontaneous fission half-life is now known reasonably well, the understanding of its half-life and its mass yield curve in terms of a theoretical model is necessary for predicting half-lives and mass yield curves of superheavy nuclei.

We examine in this article the decay of these isotopes within the context of the recently developed theory of fission [3, 4] using an interaction characterized by a thin external barrier. Such a barrier follows directly from the computation of the interaction potential using a density dependent mass formula [3, 5, 6], i.e., a mass formula which explicitly incorporates in the energy expression the appropriate variation of the nuclear density as a function of the distance from the centre of the nucleus. Such a mass formula [7–9] accounts for known masses of nuclei as good as any standard mass formula obtained from a liquid droplet of a constant density. Using a simple parametrized version of such a barrier, we have already examined various observed properties associated with the spontaneous fission of ²³⁴U, ²³⁶U, ²⁴⁰Pu and ²⁵²Cf (Refs. [10] and [11]) and found substantial agreement between the theoretical calculations and observed data. We apply this theory and use the same set of parameters as those in Refs. [10] and [11]

On a leave of absence from Physics Department, Indiana University, Bloomington, Indiana 47401, USA. This work is performed as a visiting professor under the auspices of the commission of physics, Switzerland (C.I.C.P.).

to calculate various observables associated with the spontaneous decay of ²⁵⁰Cm, ²⁵⁴Cf and ²⁵⁸Fm. We also compute their kinetic energy versus mass spectra using the theory of Terrell [12].

Yield curves are then computed in the JWKB approximation discussed in Refs. [22], [11] and [4]. These theoretical kinetic energy spectra associated with the decay of ²⁵⁰Cm, ²⁵⁴Cf and ²⁵⁸Fm by fission are drawn in the upper half of Figures 1, 2 and 3, respectively, by solid points. Dashed curves in these figures represent actual values used in computing the mass yield curve. Numerals next to solid points refer to the atomic number of the heavy fragment of a particular daughter pair used in calculating the kinetic energy. The lower half of these figures plots the theoretical percentage yield curves as a function of the mass of the heavy fragment.

We clearly see that while the dominant decay modes of ²⁵⁰Cm and ²⁵⁴Cf are non-symmetric, ²⁵⁸Fm prefers to decay primarily in the symmetric daughter pairs. This is consistent with observations.

Table I Experimental and theoretical values of spontaneous fission half-lives (referred, respectively, as $E_{\rm kin}({\rm exp})$ and $E_{\rm kin}({\rm th})$) for the fission of a number of transuranic elements. Half-lives and kinetic energies of elements marked with an (*) refer only to the decay to their respective fastest decay modes. Energies and half-lives are given in MeV and years, respectively

Element	$E_{kin}(\exp)$	$E_{ exttt{kin}}(ext{th})$	$T_{1/2}$ SFS (exp)	$T_{1/2}$ SFS(th)
²⁴⁴ Cm*	185.5 ± 5.0 [19]	185	$1.34 \cdot 10^7 [17]$	$6.80 \cdot 10^7$
²⁵⁰ Cm (set 1)	attouristations — variables Co-cost I	180.5	$2.00 \cdot 10^4 [16]$	$2.20 \cdot 10^{6}$
(set 2)		182.3		$3.10 \cdot 10^{4}$
246Cf*	195.6 ± 2.0 [20]	196	$2.00 \cdot 10^3 [21]$	$3.80 \cdot 10^3$
²⁵² Cf	186.5 ± 1.2 $\begin{bmatrix} 13 \end{bmatrix}$	186.7	85.5 [17]	86
²⁵⁴ Cf	$185.0 \pm 2.0 [14]$	188.6	$1.78 \cdot 10^{-1} [18]$	$2.40 \cdot 10^{-1}$
²⁵⁸ Fm (set 1)	$190 \sim 200 [15]$	200.7	$3.8 \cdot 10^{-11} [2]$	$3.90 \cdot 10^{-9}$
(set 2)		203.3		$5.10 \cdot 10^{-11}$

In Table I we have compared theoretically calculated total spontaneous decay half-lives and average kinetic energy with experimental data. The theoretical computations for the decay of ²⁵⁰Cm and ²⁵⁸Fm are given for two sets of kinetic energy spectrum. Set 1 corresponds to the dashed curve plotted in the upper halves of Figures 1, 2 and 3 and set 2 corresponds to adding 2 MeV extra to our computed kinetic energy for all decay modes. The mass yield curves corresponding to set 2 look essentially the same as those plotted in Figures 1, 2 and 3 using set 1.

In the table we have also reproduced our results of the decay of 252 Cf from Ref. [11] and the theoretically computed partial decay half-lives associated with the decay of 244 Cm and 246 Cf to one of their fastest decay modes (i.e., 244 Cm \rightarrow 144 Ba + 100 Zr and 246 Cf \rightarrow 144 Ce + 102 Zr) using the same set of parameters.

Thus, the theory is in a position to explain

- i) the decay of ²⁵⁸Fm predominantly to symmetric decay modes,
- ii) the variation of spontaneous fission half-lives among the even-even isotopes of Curium and Californium and
- iii) absolute total half-lives of spontaneous fission.

The reasons for the success of this simple theory in correlating so many data is evident from the discussion in Ref. [22], elucidating the relation between decay constant, kinetic energy and fission barrier.

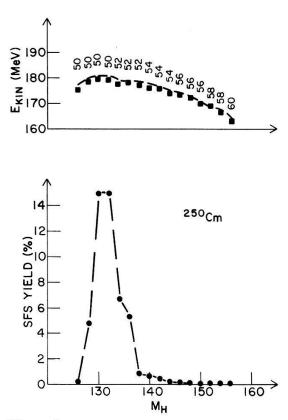


Figure 1 The kinetic energy spectrum associated with the binary spontaneous fission of $^{250}\mathrm{Cm}$. M_H refers to the heavier of the two fragments in the decay product and 'SFS yield %' means the percentage mass yield curve in the spontaneous fission.

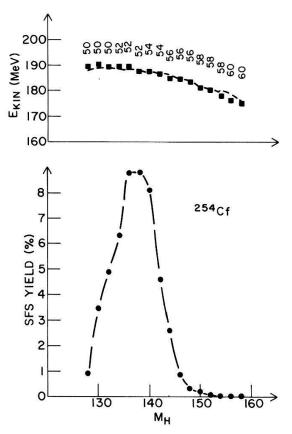
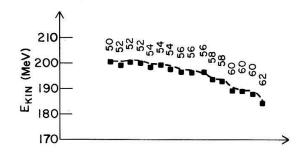


Figure 2
The same as in Figure 1 for the binary spontaneous fission of ²⁵⁴Cf.



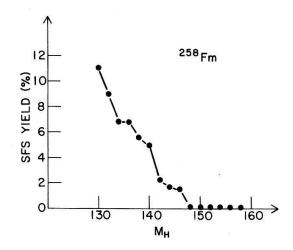


Figure 3

The same as in Figure 1 for the binary spontaneous fission of ²⁵⁸Fm. Note that the yield curve is symmetric with respect to the mass splitting.

REFERENCES

- [1] W. John, E. K. Hulet, R. W. Longheed and J. J. Wesolowski, Phys. Rev. Letters 27, 45 (1971).
- [2] E. K. HULET, J. F. WILD, R. W. LONGHEED, J. E. EVANS and B. J. QUALHEIM, Phys. Rev. Letters 26, 523 (1971).
- [3] B. BLOCK, J. W. CLARK, M. D. HIGH, R. MALMIN and F. B. MALIK, Ann. Phys (NY) 62, 464 (1971) and Bull. Am. Phys. Soc. 15, 646 (1970).
- [4] M. A. Hooshyar and F. B. Malik, Helv. Phys. Acta 45, 567 (1972).
- [5] I. REICHSTEIN, Bull. Am. Phys. Soc. 16, 516 (1971).
- [6] I. REICHSTEIN and F. B. MALIK, Phys. Letters 37B, 344 (1971).
- [7] K. A. Brueckner, J. R. Buchler, S. Jorna and R. J. Lombard, Phys. Rev. 171, 1188 (1968).
- [8] K. A. BRUECKNER, J. R. BUCHLER, R. C. CLARK and R. J. LOMBARD, Phys. Rev. 181, 1543 (1969).
- [9] R. J. LOMBARD, Ann. Phys. (N.Y.) 77, 380 (1973).
- [10] M. A. HOOSHYAR and F. B. MALIK, Phys. Letters 38B, 495 (1972).
- [11] M. A. Hooshyar and F. B. Malik, Helv. Phys. Acta (to be published).
- [12] J. TERELL, Phys. Rev. 127, 880 (1962).
- [13] H. W. SCHMITT, J. H. NEILER and F. J. WALTER, Phys. Rev. 191, 1146 (1966).
- [14] R. Brandt, S. G. Thompson, R. C. Gatti and L. Phillips, Phys. Rev. 131, 2617 (1963).
- [15] W. John, E. K. Hulet, R. W. Longheed and J. J. Wesolowski, Phys. Letters 27, 45 (1971), quotes 191 MeV but indicates that their result may lie about 10 MeV lower in some cases, e.g., at the asymmetric end of their mass distribution, their energies are about 10 MeV lower than those of J. P. Balaga, G. P. Ford, D. C. Hoffman and J. D. Knight, Phys. Rev. Letters 26, 145 (1971).
- [16] J. HUIZENGA and H. DIAMAND, Phys. Rev. 107, 1087 (1957).
- [17] E. K. Hyde, The Nuclear Properties of the Heavy Elements, Vol. III (Prentice Hall, Inc., Englewood Cliffs, NJ 1964).

- [18] L. Phillips et al., J. Inorg. Nucl. Chem. 25, 1085 (1963).
- [19] A. Smith et al., Second International Conference on Peaceful use of Atomic Energy, U.N. Publication (1959).
- [20] A. M. Friedman et al., Phys. Rev. 131, 1203 (1963).
- [21] E. K. Hulet, Phys. Rev. 89, 878 (1953).
- [22] F. B. Malik and P. C. Sabatier, Helv. Phys. Acta (to be published in 1973).